

24 janvier 2012

# Effects of massive neutrinos on the large-scale structure of the universe in standard cosmology.

*Lesgourgues, Pastor, '06 Phys Rep. 429, 307*

*Abazajian et al., arXiv: 1103.5083*

*Y. Wong, talk in PONT, 2011*

# The neutrino background

# Cosmic neutrino background...

- Prediction of the standard hot big bang.
- Process of decoupling fixed by **weak interactions**.

– **Temperature** today:  $T_{\nu,0} = \left(\frac{4}{11}\right)^{1/3} T_{\text{CMB},0} = 1.95 \text{ K}$

– **Number density** per flavour:  $n_{\nu,0} = \frac{6}{4} \frac{\zeta(3)}{\pi^2} T_{\nu,0}^3 = 112 \text{ cm}^{-3}$

– **Energy density** per flavour:  $\Omega_{\nu} h^2 = \frac{m_{\nu}}{93 \text{ eV}}$

Neutrinos can be a **significant component** of the total **dark matter** content.

If  $m_{\nu} > 1 \text{ meV}$

Minimum amount of  
neutrino dark matter

$\min \sum m_{\nu} \sim 0.05 \text{ eV} \rightarrow \min \Omega_{\nu} \sim 0.1 \%$

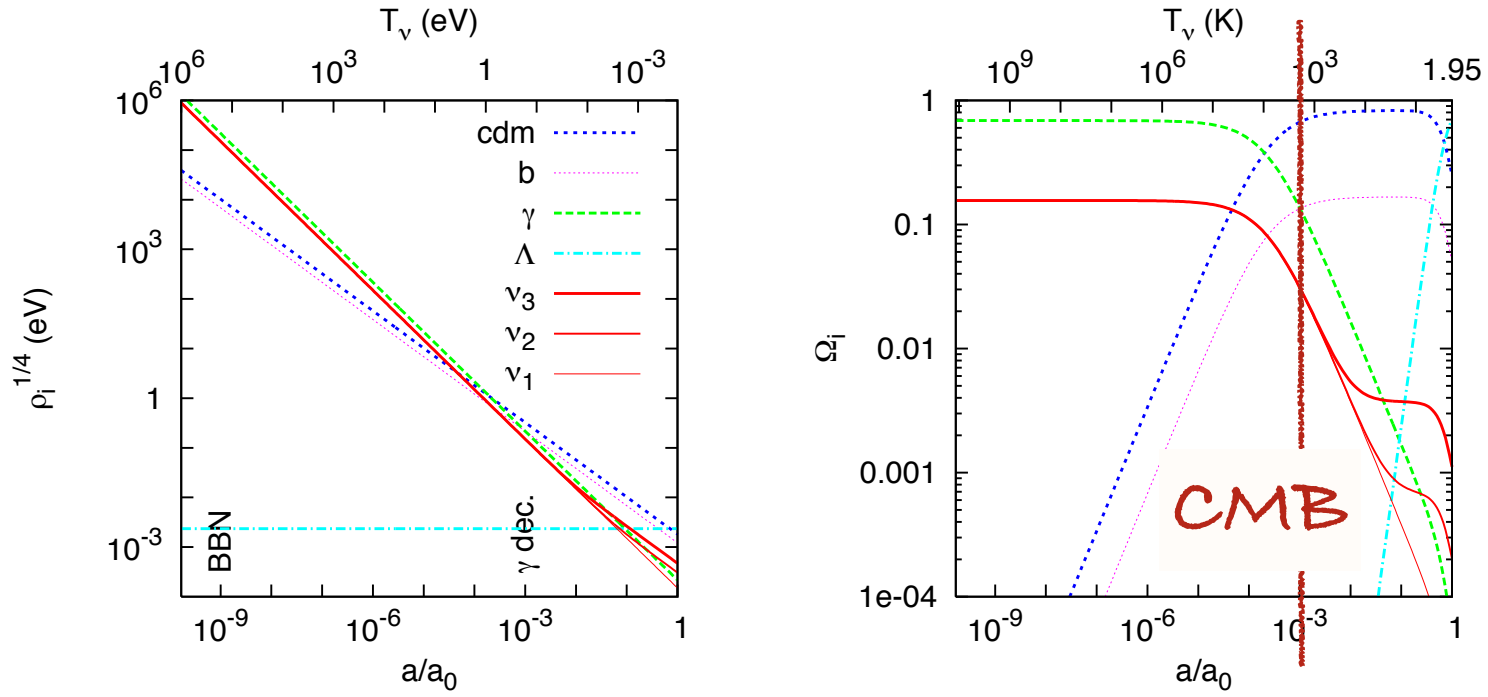


Fig. 5. Evolution of the background densities from the time when  $T_\nu = 1$  MeV (soon after neutrino decoupling) until now, for each component of a flat  $\Lambda$ MDM model with  $h = 0.7$  and current density fractions  $\Omega_\Lambda = 0.70$ ,  $\Omega_b = 0.05$ ,  $\Omega_\nu = 0.0013$  and  $\Omega_{\text{cdm}} = 1 - \Omega_\Lambda - \Omega_b - \Omega_\nu$ . The three neutrino masses are distributed according to the Normal Hierarchy scheme (see Sec. 2) with  $m_1 = 0$ ,  $m_2 = 0.009$  eV and  $m_3 = 0.05$  eV. On the left plot we show the densities to the power  $1/4$  (in eV units) as a function of the scale factor. On the right plot, we display the evolution of the density fractions (i.e., the densities in units of the critical density). We also show on the top axis the neutrino temperature (on the left in eV, and on the right in Kelvin units). The density of the neutrino mass states  $\nu_2$  and  $\nu_3$  is clearly enhanced once they become non-relativistic. On the left plot, we also display the characteristic times for the end of BBN and for photon decoupling or recombination.

# Effects on LSS growth: an overview

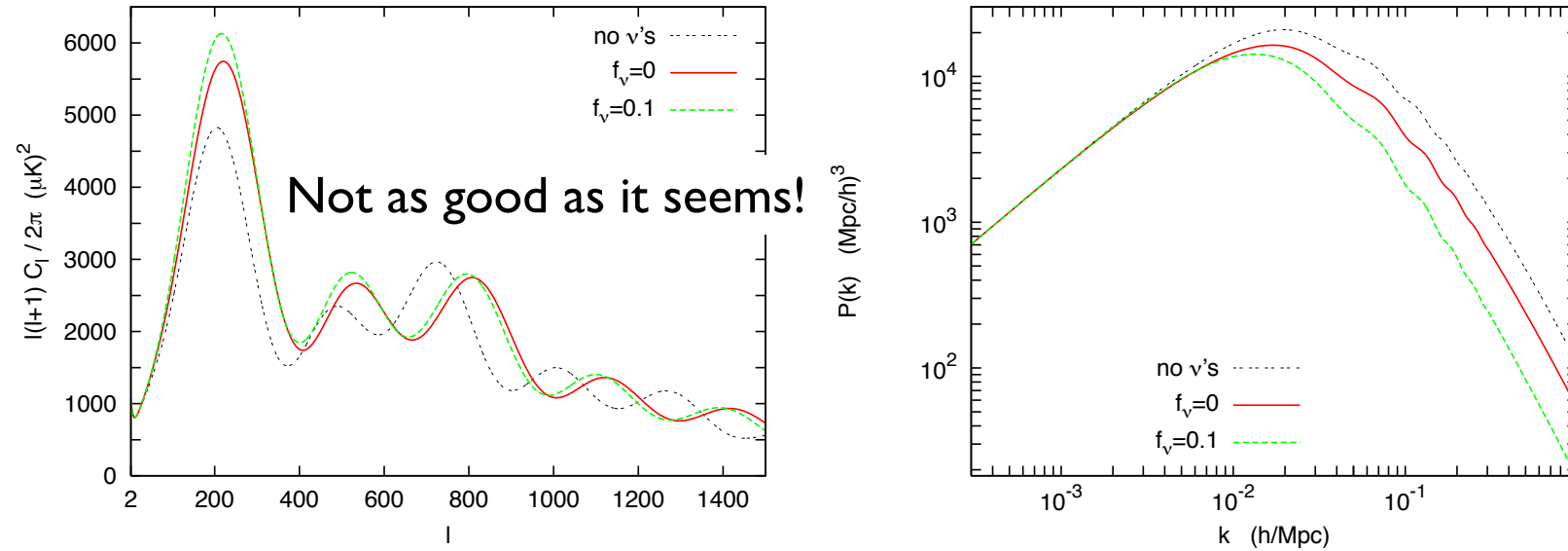


Fig. 14. CMB temperature anisotropy spectrum  $C_l^T$  and matter power spectrum  $P(k)$  for three models: the neutrinoless  $\Lambda$ CDM model of section 4.4.6, a more realistic  $\Lambda$ CDM model with three massless neutrinos ( $f_\nu \simeq 0$ ), and finally a  $\Lambda$ MDM model with three massive degenerate neutrinos and a total density fraction  $f_\nu = 0.1$ . In all models, the values of  $(\omega_b, \omega_m, \Omega_\Lambda, A_s, n, \tau)$  have been kept fixed.

# The linear growth rate with neutrinos

For dark matter

$$\dot{\delta}_c(\mathbf{k}, \tau) + \theta_c(\mathbf{k}, \tau) = 0$$

$$\dot{\theta}_c(\mathbf{k}, \tau) + \mathcal{H}\theta_c(\mathbf{k}, \tau) + \frac{3}{2}\mathcal{H}^2\delta_{\text{total}}(\mathbf{k}, \tau) = 0$$

For neutrinos (effective equations)

$$\dot{\delta}_\nu(\mathbf{k}, \tau) + \theta_\nu(\mathbf{k}, \tau) = 0$$

$$\dot{\theta}_\nu(\mathbf{k}, \tau) + \mathcal{H}\theta_\nu(\mathbf{k}, \tau) + \frac{3}{2}\mathcal{H}^2\delta_{\text{total}}(\mathbf{k}, \tau) - k^2 c_s^2(\tau)\delta_\nu(\mathbf{k}, \tau) = 0$$

... a free-streaming scale

Within the freestreaming  
"horizon"

$$\ddot{\delta}_{\text{cdm}} + \frac{2}{\tau} \dot{\delta}_{\text{cdm}} - \frac{6}{\tau^2} (1 - f_\nu) \delta_{\text{cdm}} = 0$$

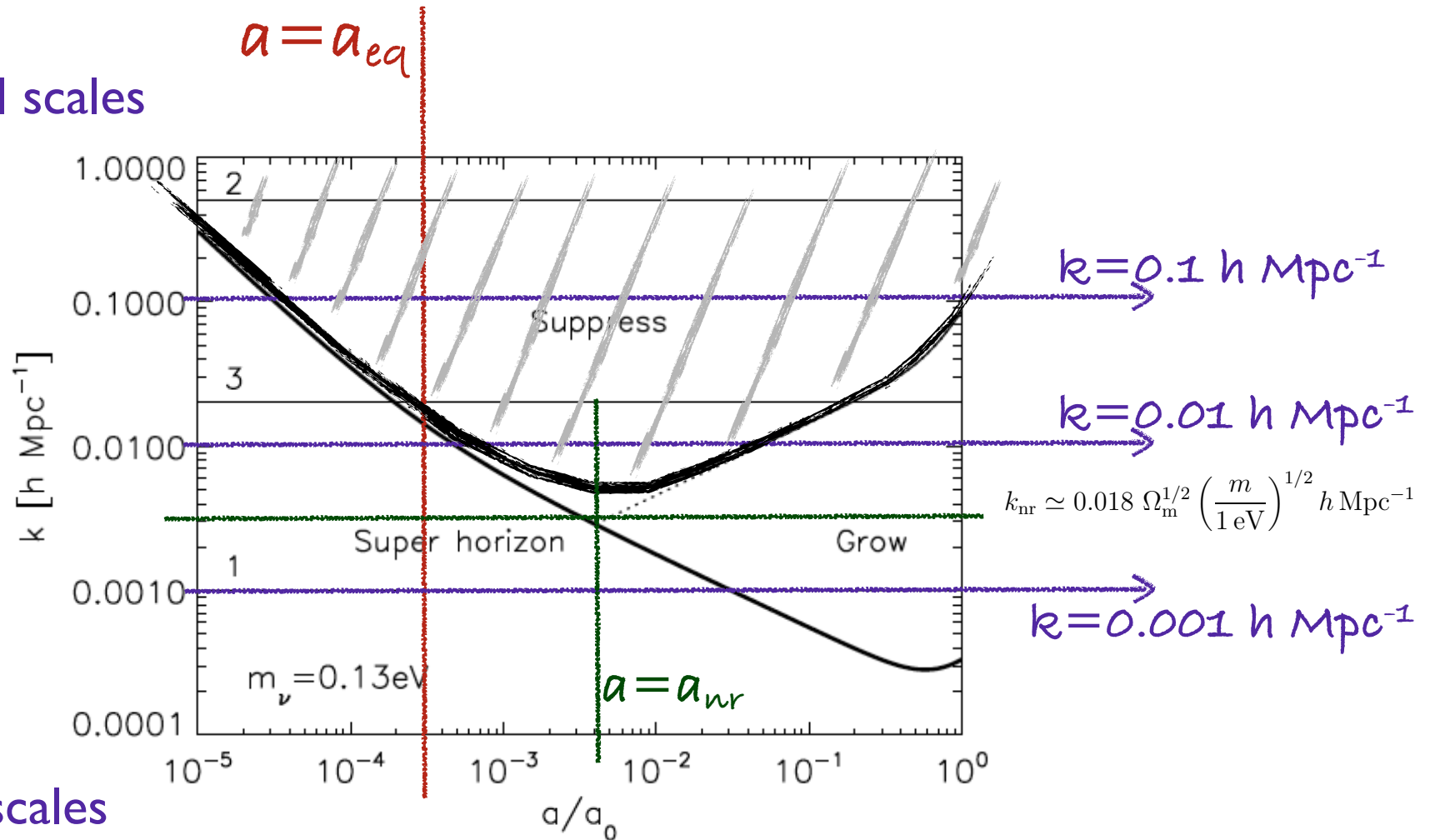
$$f_\nu \equiv \frac{\rho_\nu}{(\rho_{\text{cdm}} + \rho_{\text{b}} + \rho_\nu)} = \frac{\Omega_\nu}{\Omega_{\text{m}}} \quad f_\nu \rightarrow 0.0035$$

$$\delta_{\text{cdm}} \propto a^{(-1 + \sqrt{1 + 24(1 - f_\nu)})/4}$$

$$\approx 1 - 0.6f_\nu$$

# Free-Streaming horizon

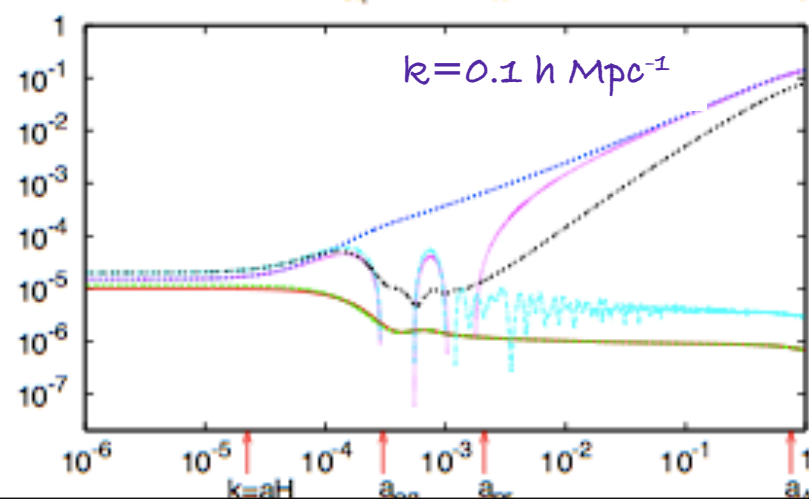
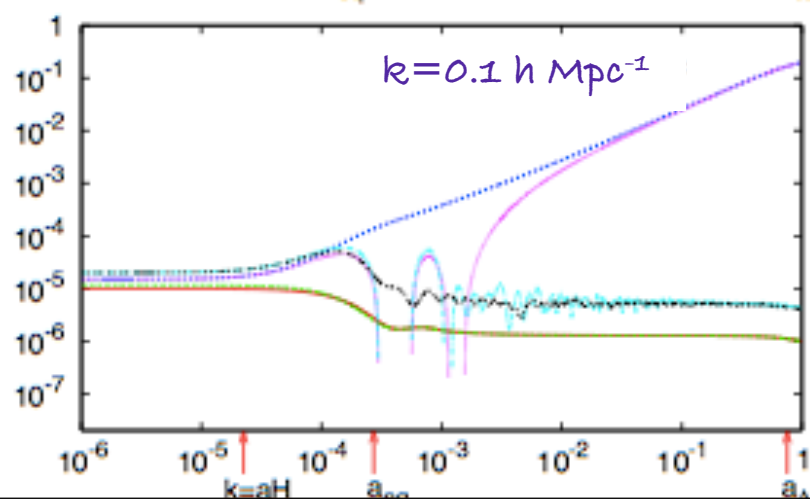
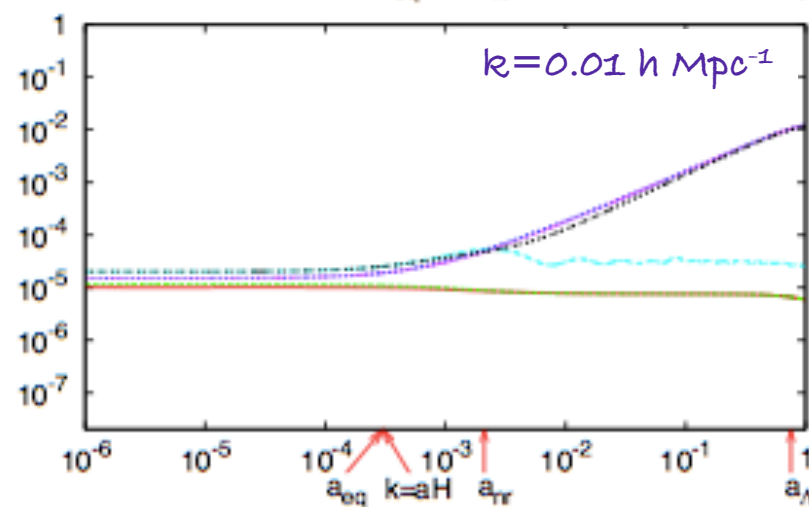
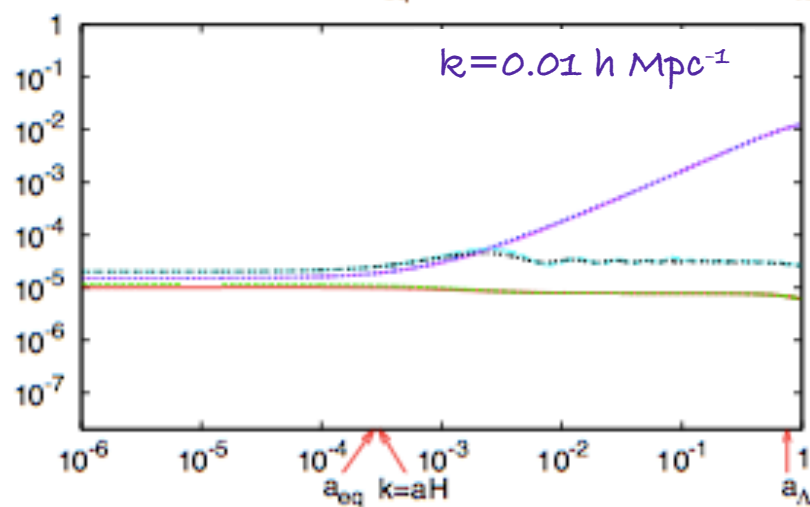
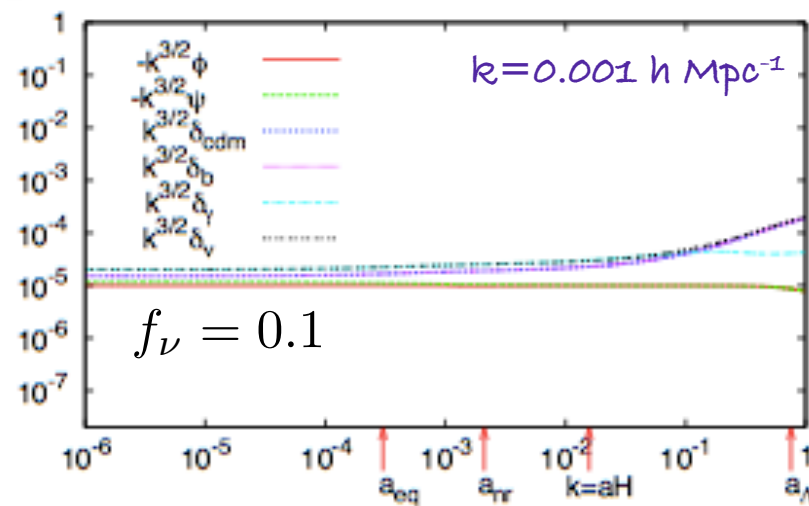
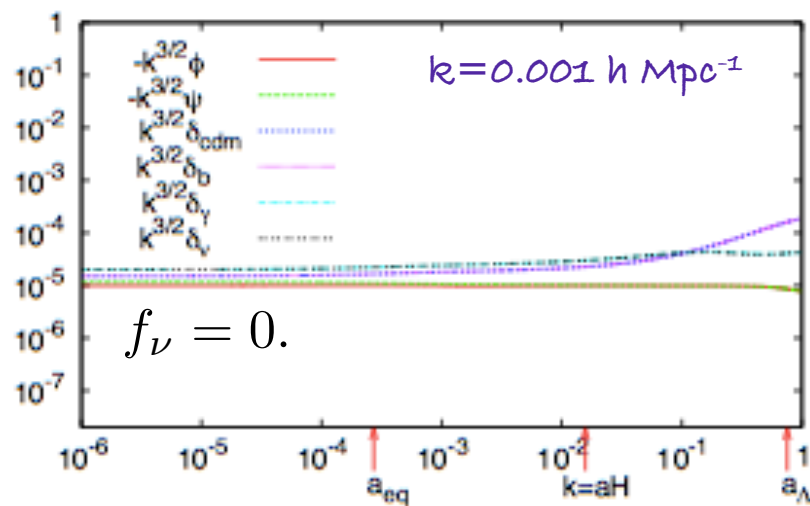
small scales



large scales

Regions where the growth of neutrino density contrast (for each species) is suppressed  $\delta_{\text{cdm}} \propto a^{1-.6f_\nu}$





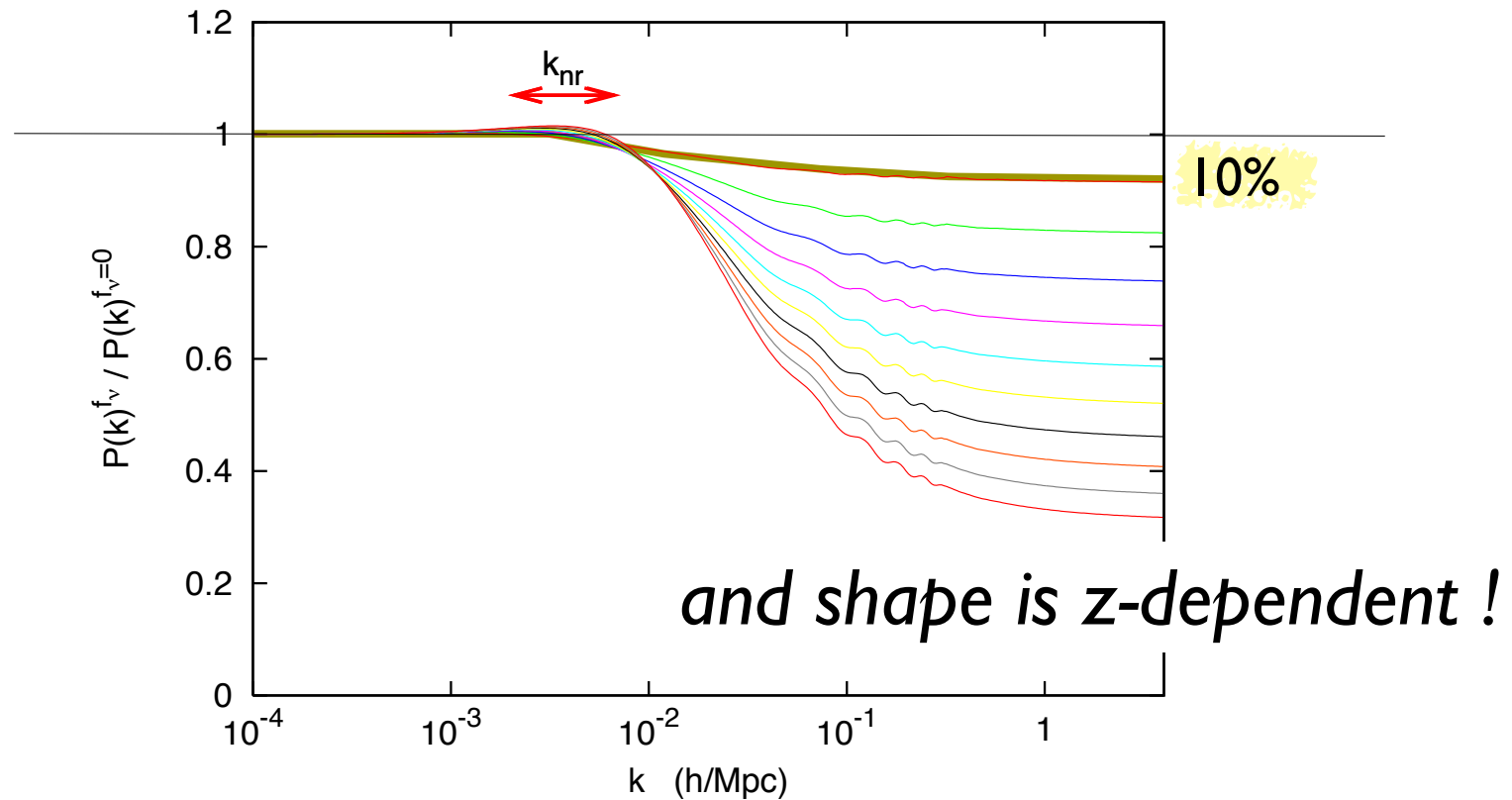
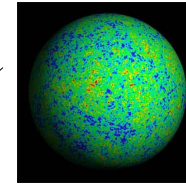
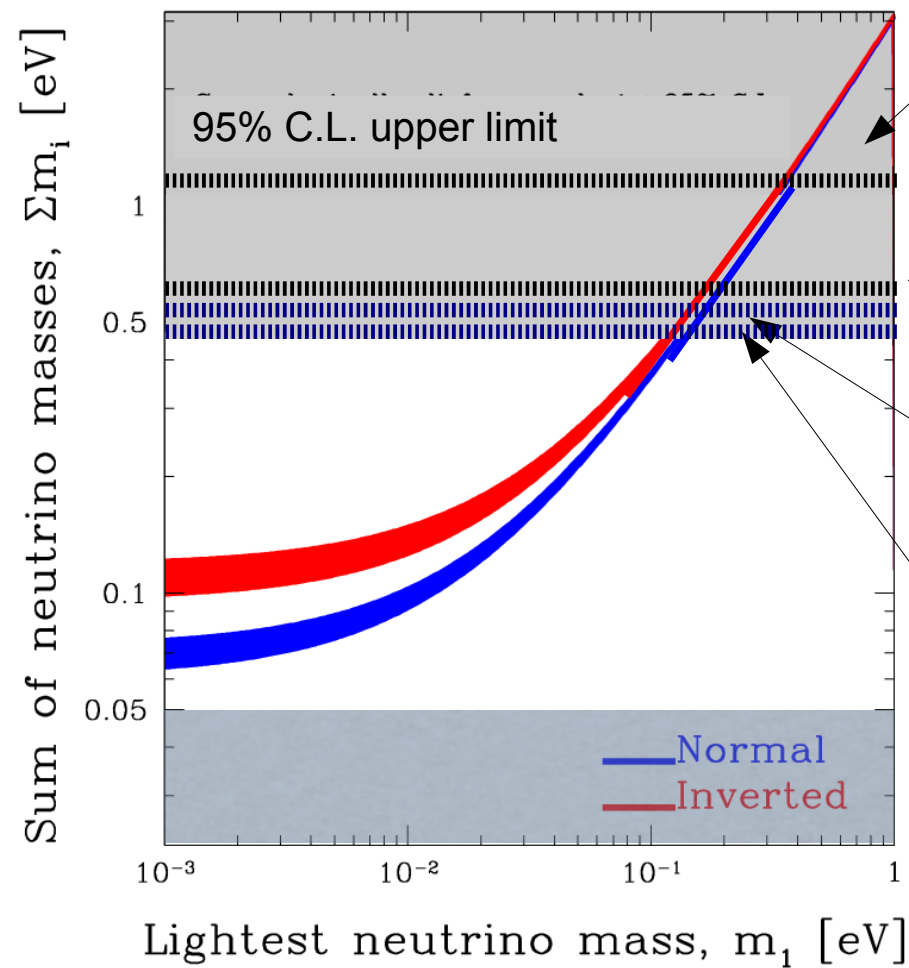
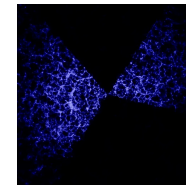


Fig. 13. Ratio of the matter power spectrum including three degenerate massive neutrinos with density fraction  $f_\nu$  to that with three massless neutrinos. The parameters  $(\omega_m, \Omega_\Lambda) = (0.147, 0.70)$  are kept fixed, and from top to bottom the curves correspond to  $f_\nu = 0.01, 0.02, 0.03, \dots, 0.10$ . The individual masses  $m_\nu$  range from 0.046 eV to 0.46 eV, and the scale  $k_{nr}$  from  $2.1 \times 10^{-3} h \text{ Mpc}^{-1}$  to  $6.7 \times 10^{-3} h \text{ Mpc}^{-1}$  as shown on the top of the figure.  $k_{eq}$  is approximately equal to  $1.5 \times 10^{-2} h \text{ Mpc}^{-1}$ .

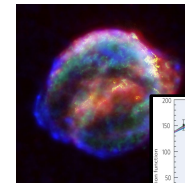
# Present status...



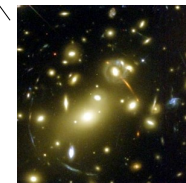
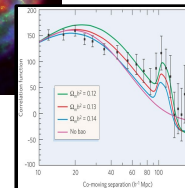
WMAP7 only\*  
Komatsu et al. 2010



WMAP7+SDSS-HPS \*  
Hannestad, Mirizzi, Raffelt  
& Y<sup>3</sup>W 2010



WMAP5+SDSS-HPS  
+SN+HST  
Reid et al. 2009  
(extended models)



WMAP5+Weak lensing\*  
Tereno et al. 2008  
Ichiki et al. 2008

\*  $\Lambda$ CDM+ $m_\nu$

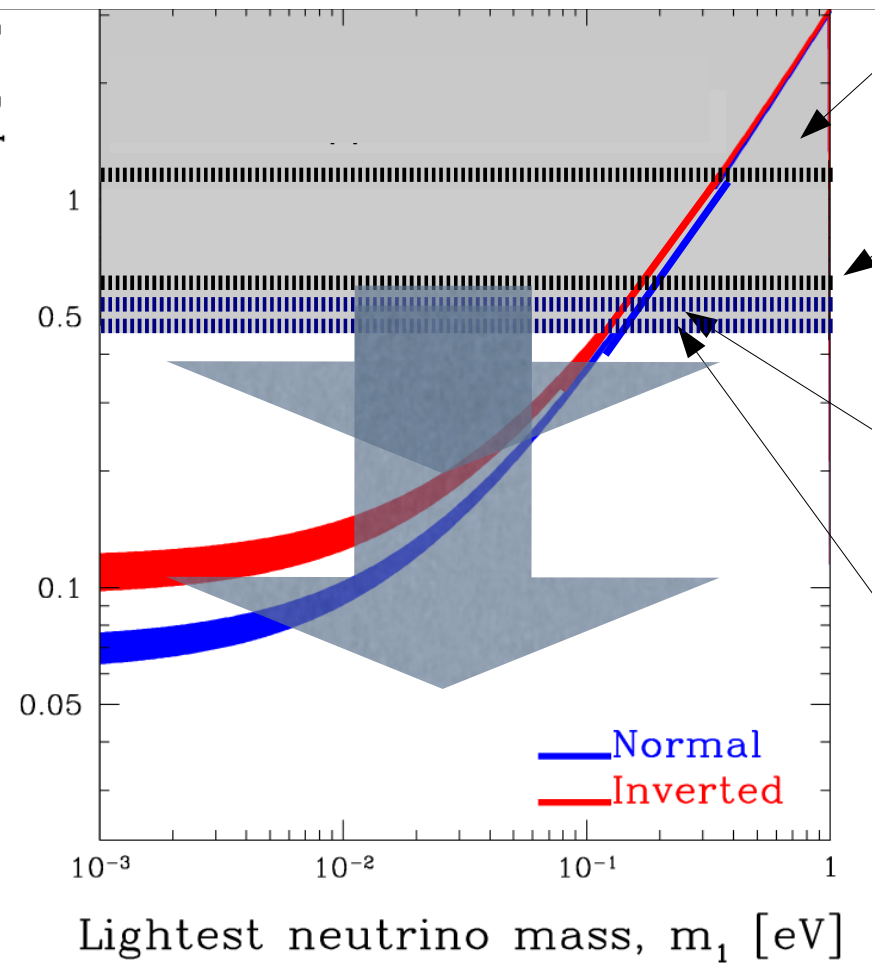
# A bit of forecast

Probe	Current $\sum m_\nu$ (eV)	Forecast $\sum m_\nu$ (eV)	Key Systematics
CMB Primordial	1.3	0.6	Recombination
CMB Primordial + Distance	0.58	0.35	Distance measurements
Lensing of CMB	$\infty$	0.2 – 0.05	NG of anisotropies
Galaxy Distribution	0.6	0.1	Nonlinearities
Lensing of Galaxies	0.6	0.07	Baryons, Metric redshifts
Lyman $\alpha$	0.2	0.1	Bias, Metals, QSO continuum
21 cm	$\infty$	0.1 – 0.006	Foregrounds, Astrophysical modeling
Galaxy Clusters	0.3	0.1	Mass Function, Mass Calibration

Planck -- CMBPol

$k=0.001-1\text{Mpc}^{-1}$  for weak lensing !

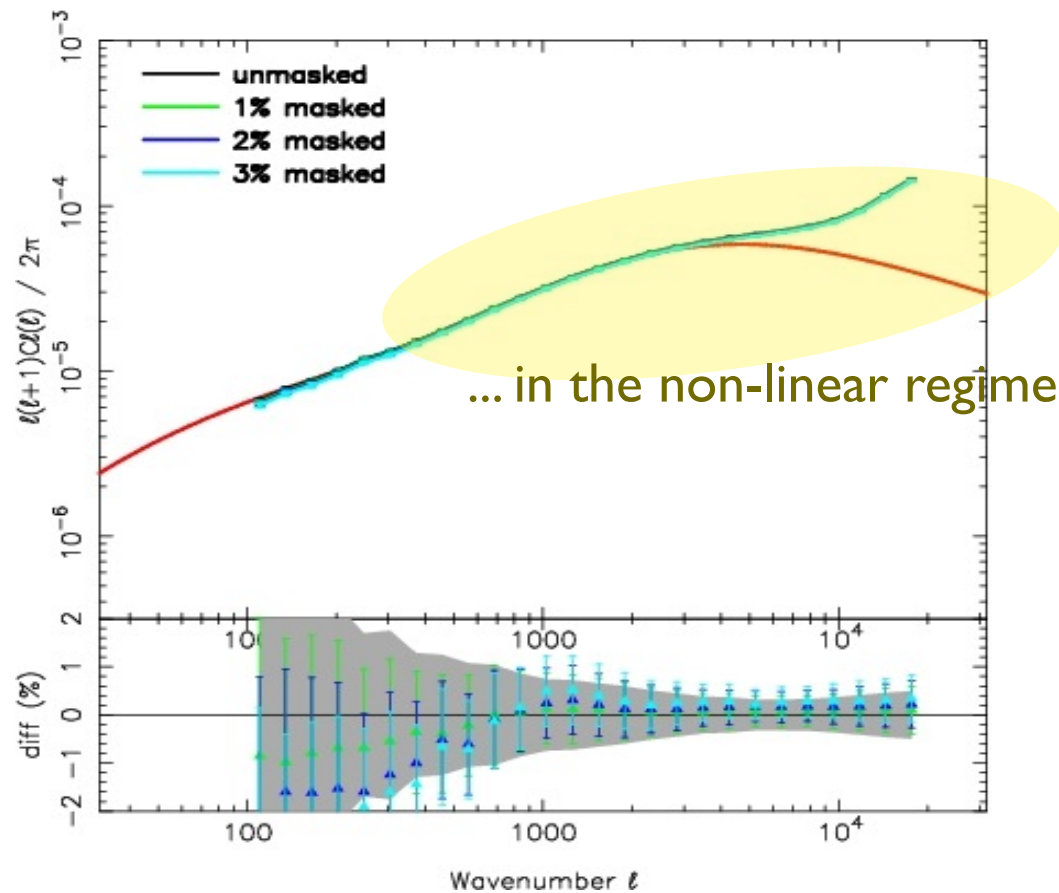
Sum of neutrino masses,  $\Sigma m_i$  [eV]



in combination with WMAP; 95% upper limits Abazajian et al. 1103.5083

a critical reading: beyond  
the linear theory...

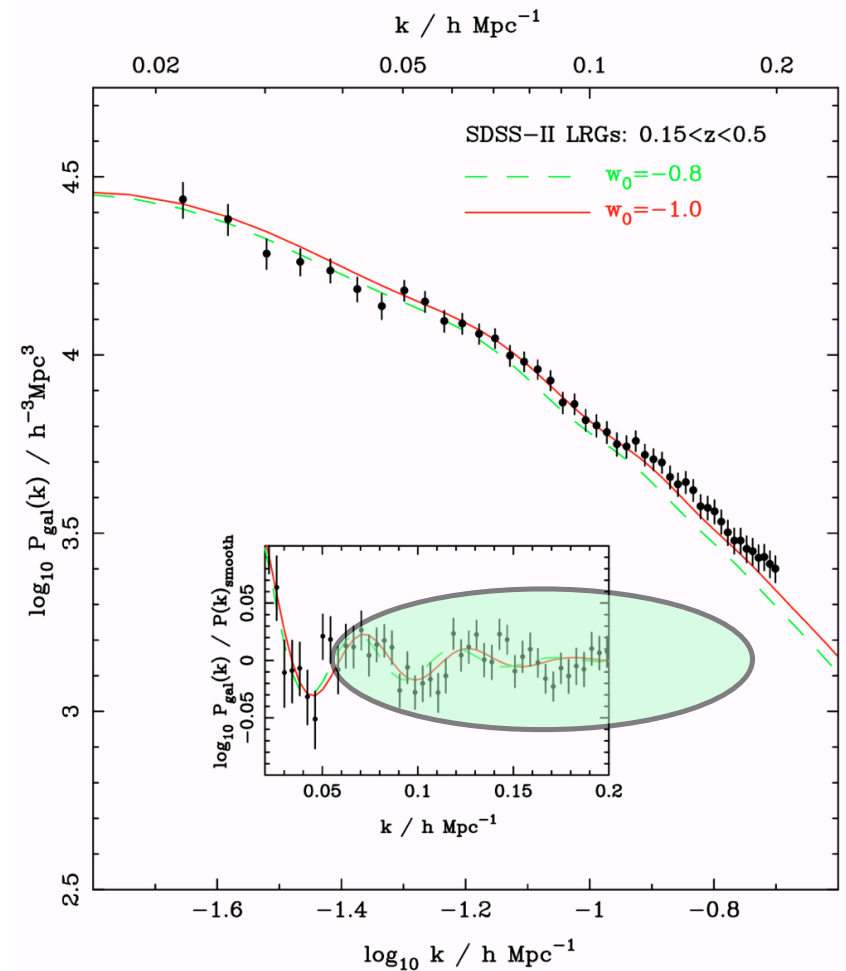
# Cosmic shear



... in the non-linear regime

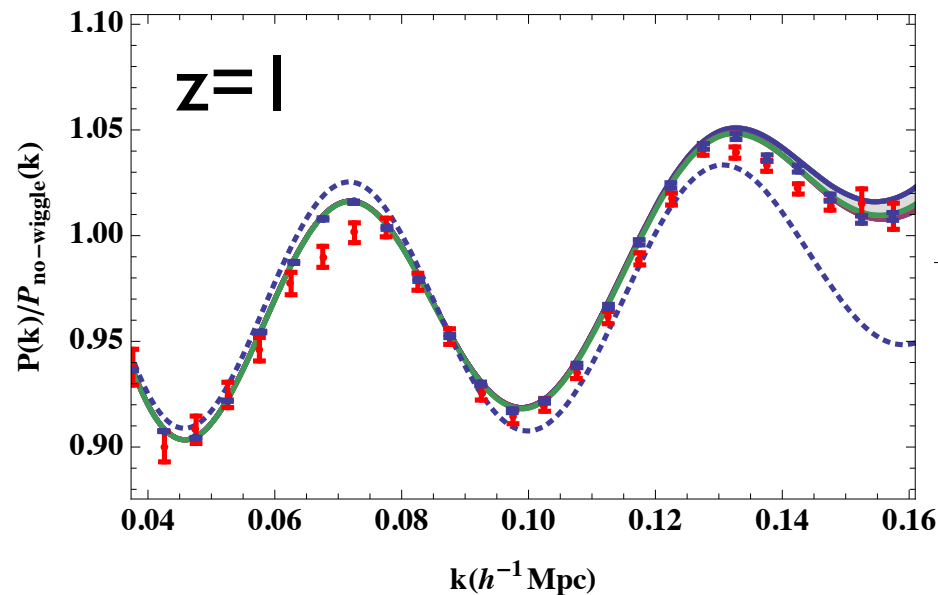
From the Euclid red book

# BAO

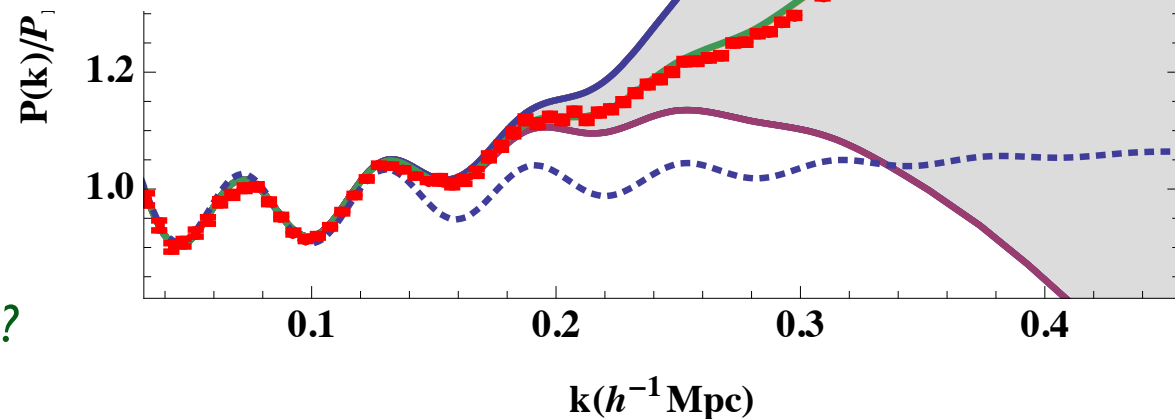


# State of the art for Perturbation Theory calculations

$$\begin{aligned}\dot{\delta}_c(\mathbf{k}, \tau) + \theta_c(\mathbf{k}, \tau) &= \alpha(\mathbf{k}_1, \mathbf{k}_2) \theta(\mathbf{k}_1) \delta(\mathbf{k}_2) \\ \dot{\theta}_c(\mathbf{k}, \tau) + \mathcal{H} \theta_c(\mathbf{k}, \tau) + \frac{3}{2} \mathcal{H}^2 \delta_{\text{total}}(\mathbf{k}, \tau) &= \beta(\mathbf{k}_1, \mathbf{k}_2) \theta(\mathbf{k}_1) \theta(\mathbf{k}_2)\end{aligned}$$



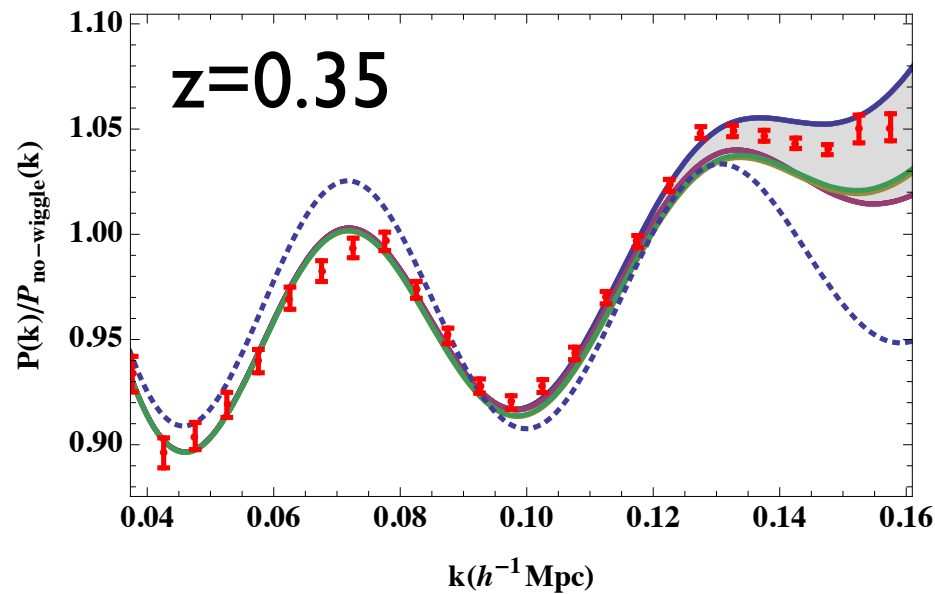
*2-loop results(!) for  $\Lambda$ -CDM*



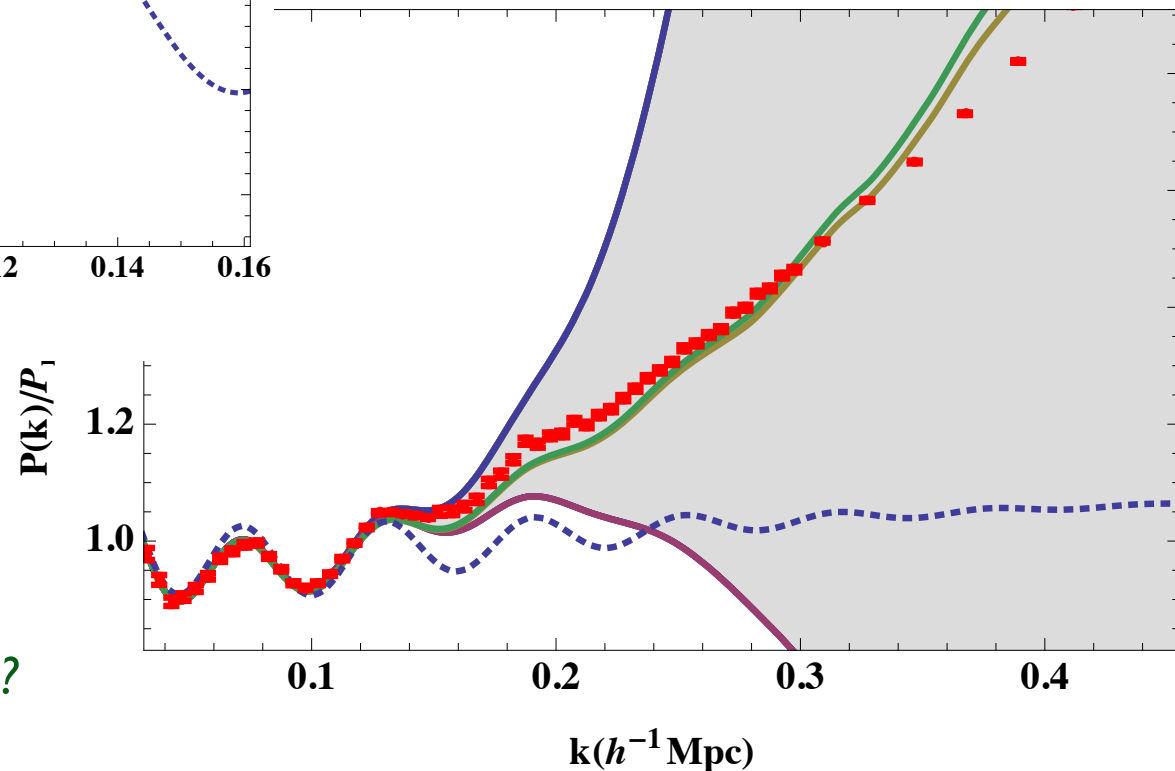
FB, A. Taruya, T. Nishimichi, '??'

# State of the art for Perturbation Theory calculations

$$\begin{aligned}\dot{\delta}_c(\mathbf{k}, \tau) + \theta_c(\mathbf{k}, \tau) &= \alpha(\mathbf{k}_1, \mathbf{k}_2) \theta(\mathbf{k}_1) \delta(\mathbf{k}_2) \\ \dot{\theta}_c(\mathbf{k}, \tau) + \mathcal{H} \theta_c(\mathbf{k}, \tau) + \frac{3}{2} \mathcal{H}^2 \delta_{\text{total}}(\mathbf{k}, \tau) &= \beta(\mathbf{k}_1, \mathbf{k}_2) \theta(\mathbf{k}_1) \theta(\mathbf{k}_2)\end{aligned}$$



*2-loop results(!) for  $\Lambda$ -CDM*



FB, A.Taruya, T. Nishimichi, '??'



# Conclusions

- Constraints from linear regime calculations (CMB) are solid;
- local LSS observations offer a potential for important discoveries - provided theoretical constraints are well controlled (NL evolution);
- z-evolution is key to separate from other parameters;